

Spin-dependent transport in layered magnetic metals

$\mu\Phi$

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Summary:

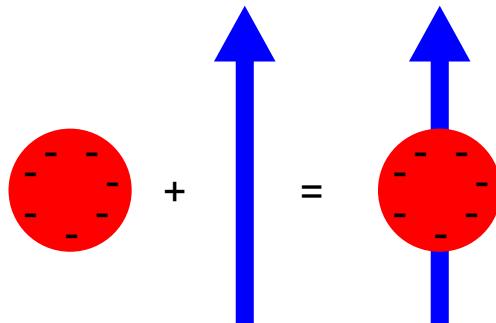
- introduction: what is spin-electronics
- giant magnetoresistance (GMR)
- tunneling magnetoresistance (TMR)
- hot-electron spin-transistor
- magnetization switching due to spin-injection
- *ab initio* calculations of perpendicular current GMR
- domain wall magnetoresistance
- theory of TMR

**introduction:
what is spin-electronics ?**

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what is an electron ?

particle with **negative electric charge** $q = -e$
and **spin 1/2 (magnetic moment $m = \mu_B$)**

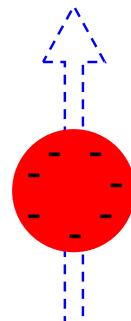


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electron as seen by an **electronician**:

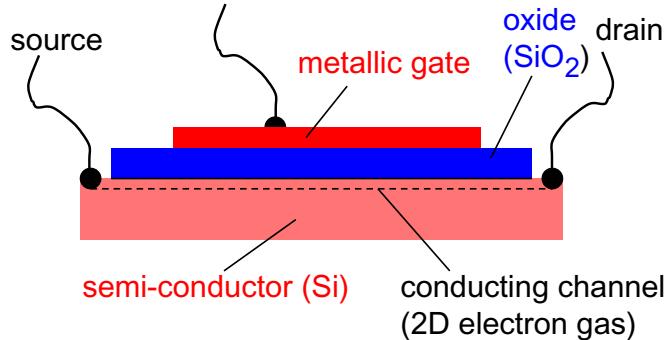
electronics = manipulation of electrons
by using their **charge** for storage and
processing of information

the **spin** is (almost) completely neglected



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principal electronic device: **MOSFET**

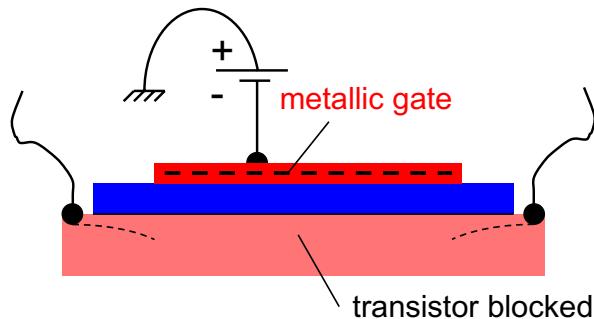


application:

- logic gates
- random access memory

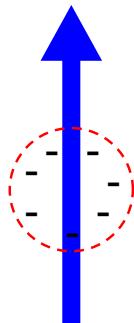
inconvenients:

- volatility of the information
- energy consumption
- limited density of information



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electron as seen by a **mystic**:



purpose of magnetism: develop materials in which the electron spins tend to align parallel to each other (**magnets**)

the **charge** of the electrons plays a secondary role

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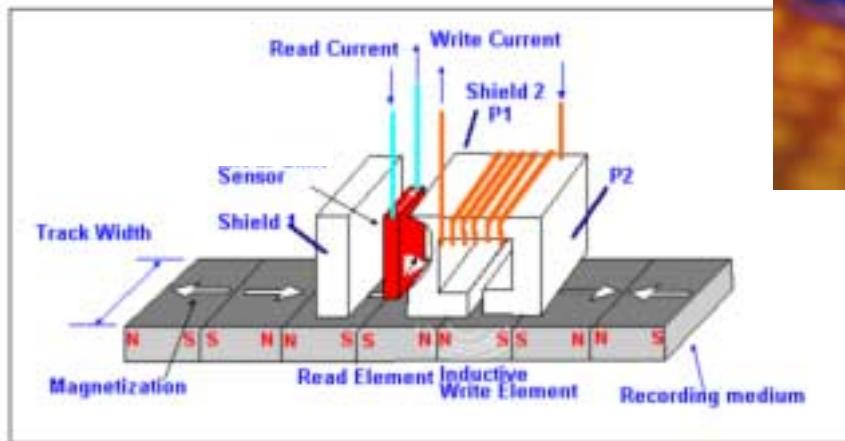
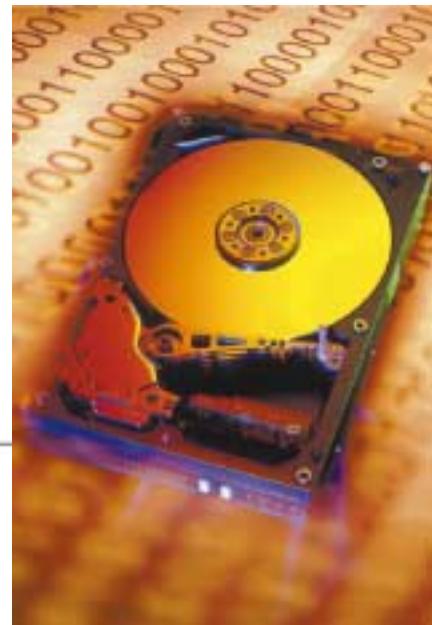
application: mass storage of information
(magnetic disks and tapes)

advantages:

- non-volatility
- high storage density
- no energy consumption

inconvenients:

- mechanical access to information



Purpose of spin-electronics: *“Teaching electrons new tricks”*

combine **electronics** and **magnetism** in order to make new devices
in which both the **charge** and the **spin** of the electron play an active role

new fundamental physical questions

new phenomena

new devices and applications

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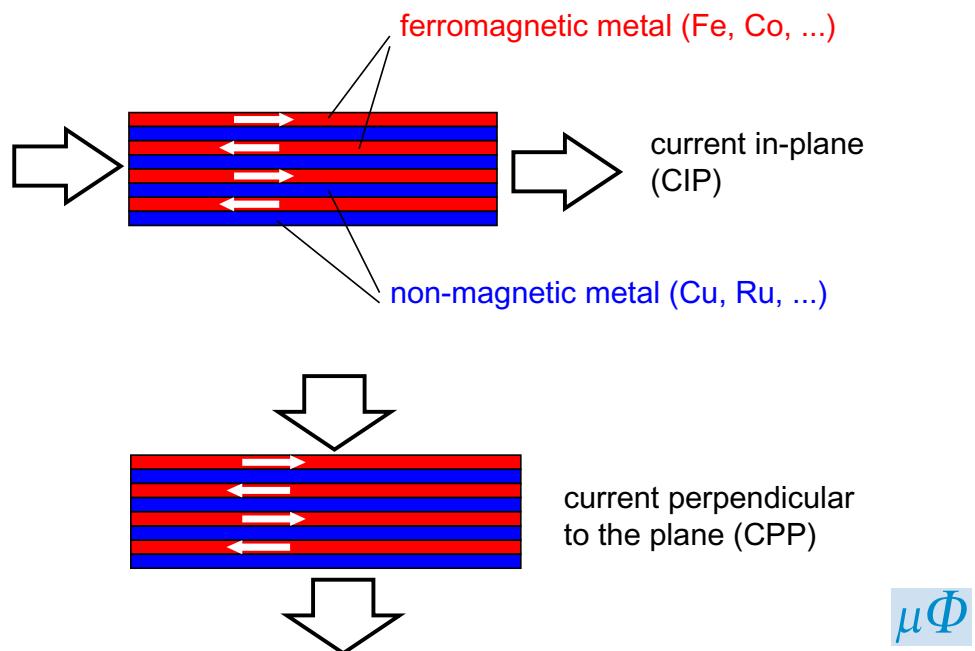
giant-magnetoresistance

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Giant magneto-resistance (GMR)

Baibich *et al.*, PRL **61**, 2472 (1988)

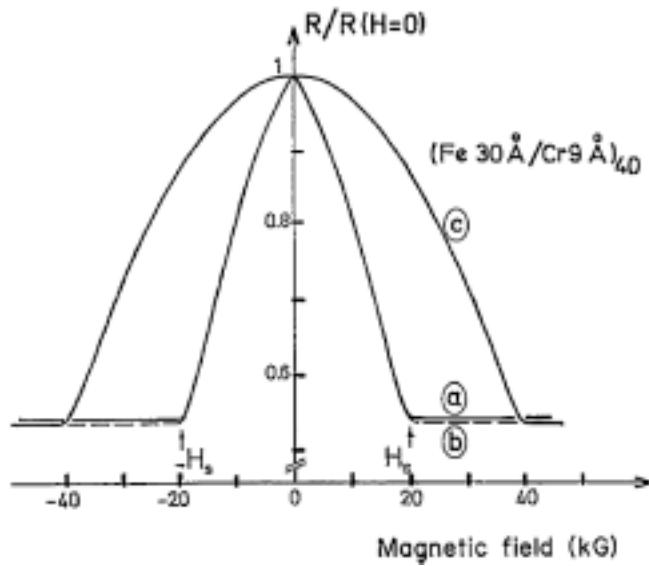
Binasch *et al.*, PRB **39**, 4828 (1989)



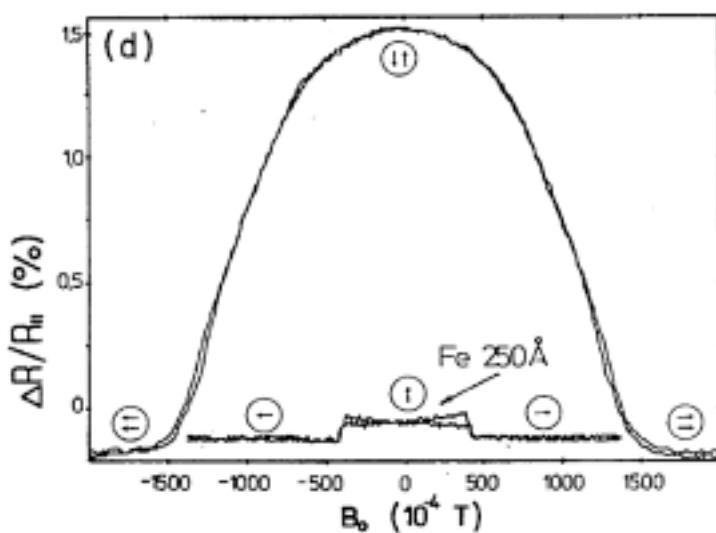
Giant Magnetoresistance of (001) Fe/(001) Cr Magnetic Superlattices

M. N. Baibich,^(a) J. M. Broto, A. Fert, F. Nguyen Van Dau, and F. Petroff
Laboratoire de Physique des Solides, Université Paris-Sud, F-91405 Orsay, France

P. Eitenne, G. Creuzet, A. Friederich, and J. Chazelas
Laboratoire Central de Recherches, Thomson CSF, B.P. 10, F-91401 Orsay, France
 (Received 24 August 1988)

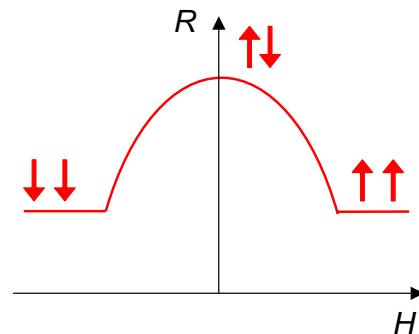

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Enhanced magnetoresistance in layered magnetic structures with antiferromagnetic interlayer exchange

G. Binasch, P. Grünberg, F. Saurenbach, and W. Zinn
Institut für Festkörperforschung, Kernforschungsanlage Jülich G.m.b.H., Postfach 1913, D-5170 Jülich, West Germany
 (Received 31 May 1988; revised manuscript received 12 December 1988)


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definition conventions for the magnetoresistance ratio

$$A = \frac{R_{AP} - R_P}{R_P} \quad \text{"optimistic" definition}$$

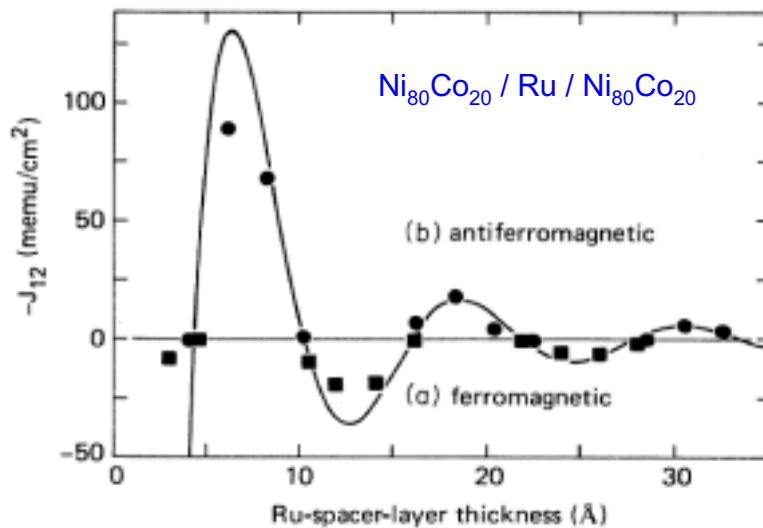


$$A = \frac{R_{AP} - R_P}{R_{AP}} \quad \text{"pessimistic" definition}$$

$$A = \frac{R_{AP} - R_P}{R_{AP} + R_P} \quad \text{reasonable definition}$$

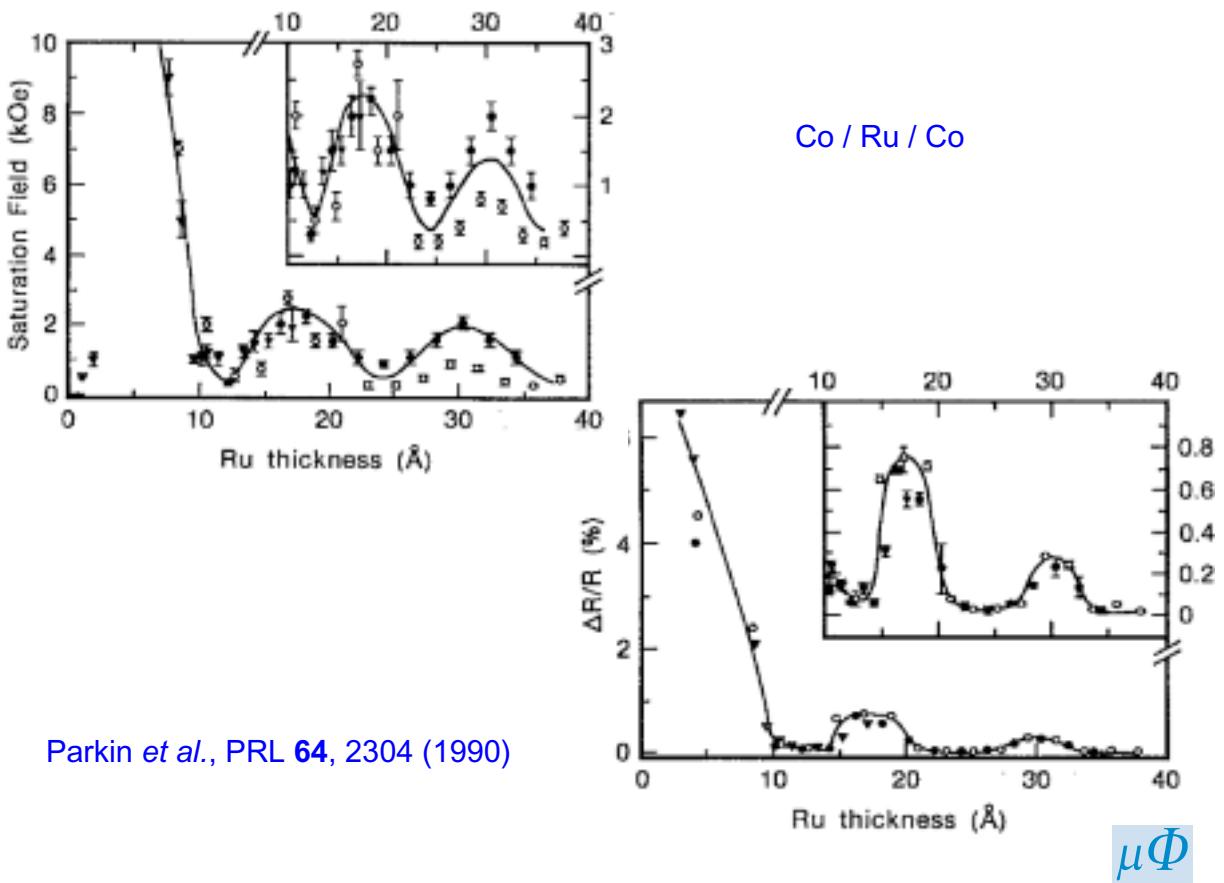
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interlayer exchange coupling

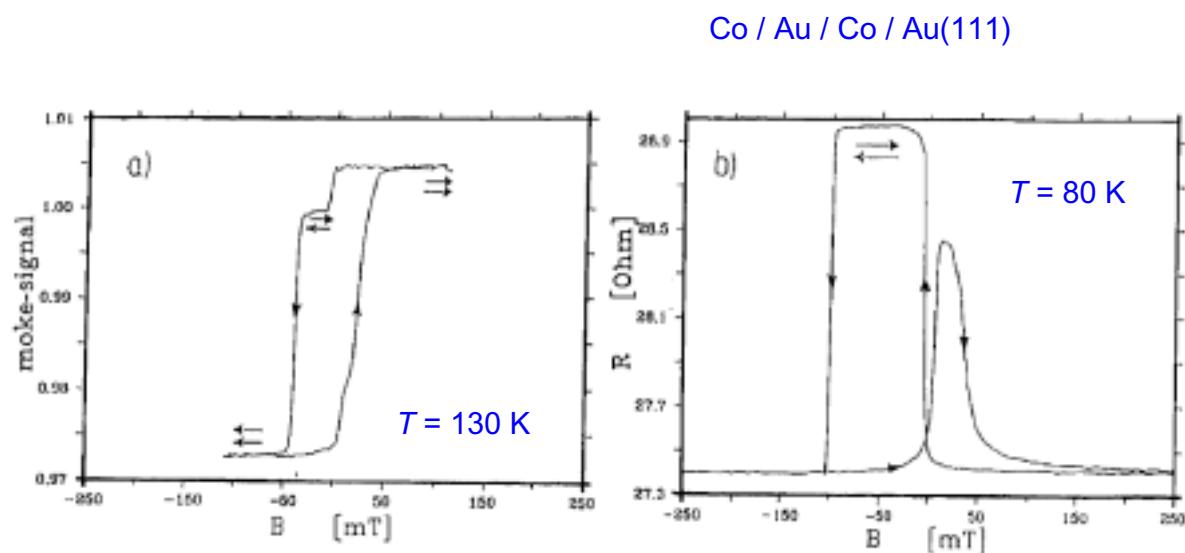


Parkin and Mauri, PRB 44, 7131 (1991)

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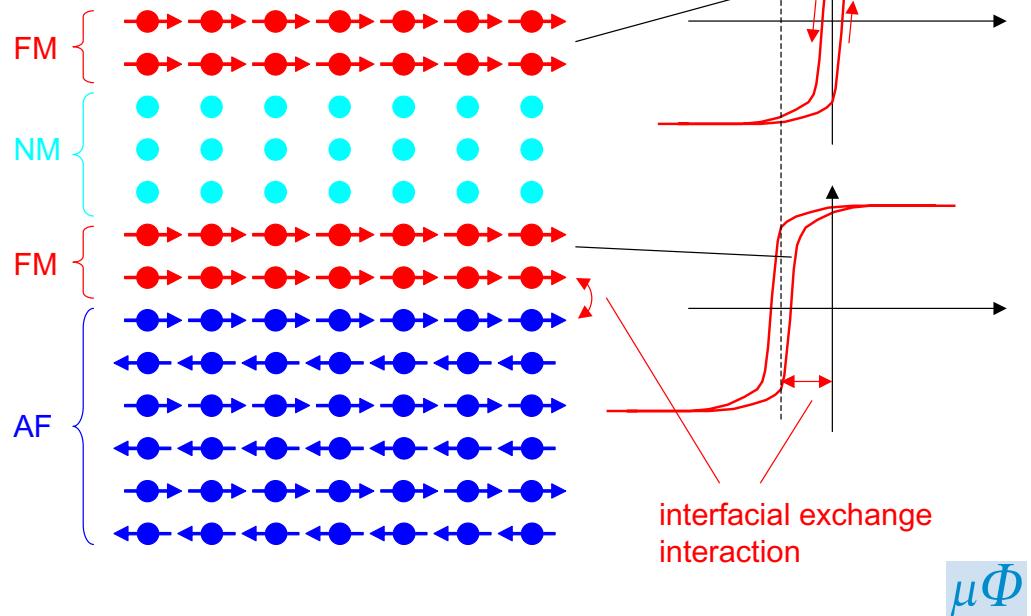
GMR without interlayer exchange coupling



Barnas *et al.*, Vacuum 41, 1241 (1990)

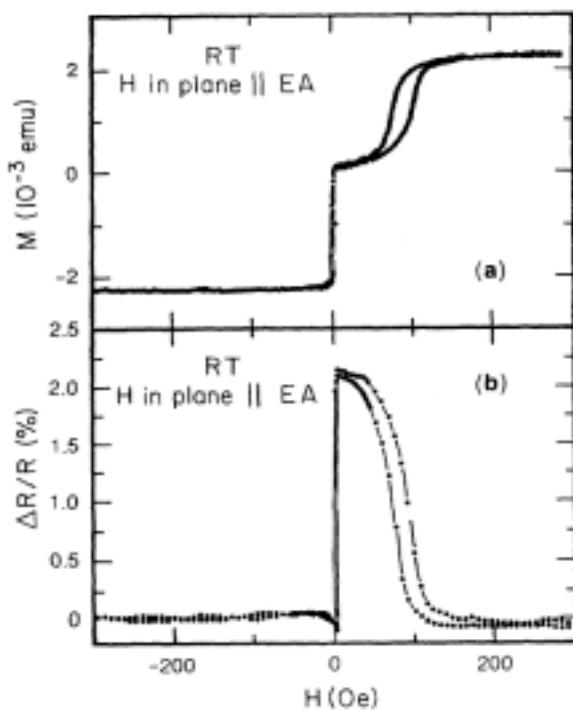
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Exchange biasing to an antiferromagnet

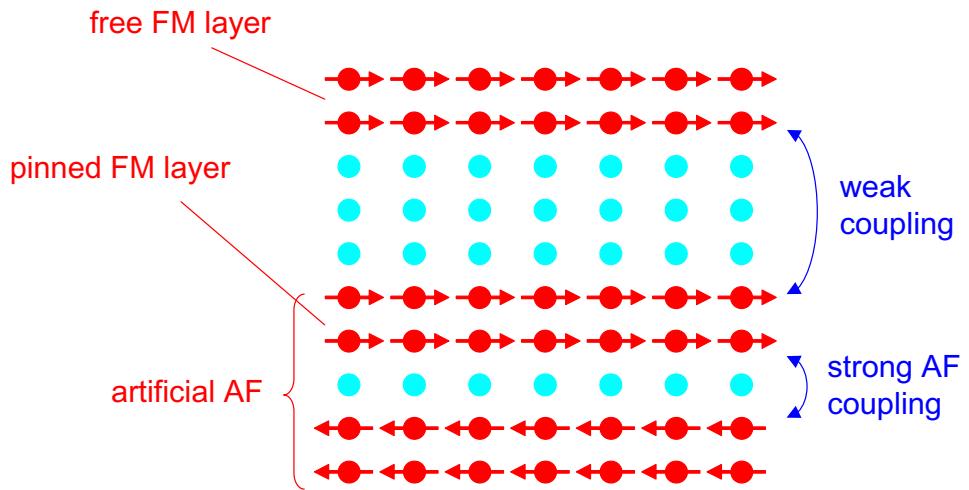


FeNi Cu FeNi FeMn

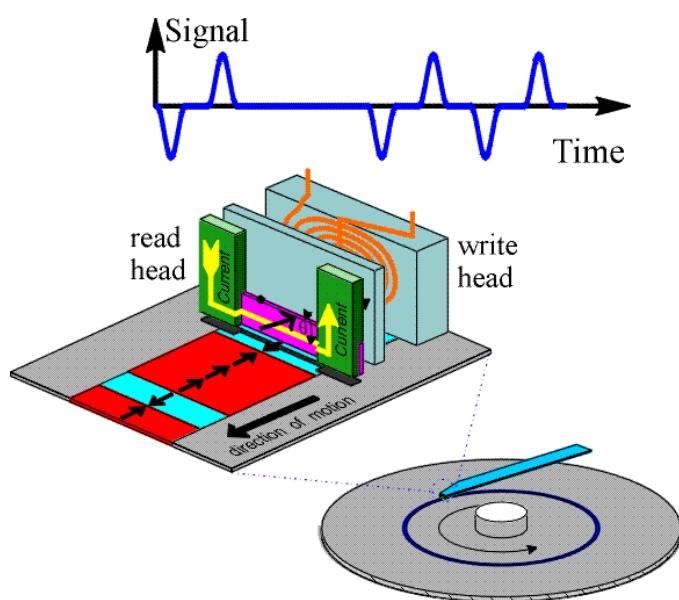
A schematic diagram showing the stack sequence of the exchange biased structure. It consists of four layers: FeNi, Cu, FeNi, and FeMn, stacked vertically. The FeNi and FeMn layers are represented by red dots, the Cu layer by cyan dots, and the second FeNi layer by blue dots.



Biasing by artificial antiferromagnet

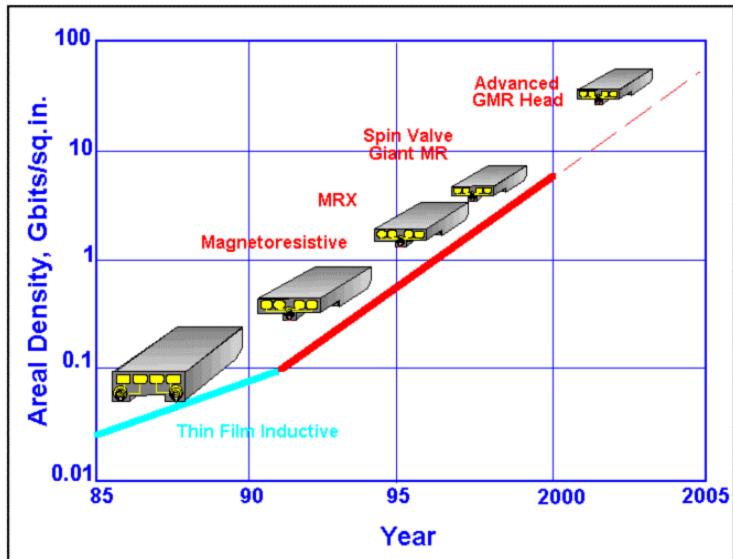


Applications of GMR: reading head for magnetic disks



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Evolution of magnetic storage density

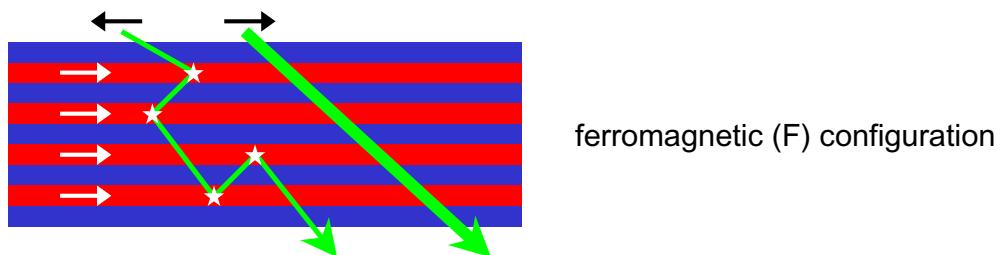


1 GByte drive

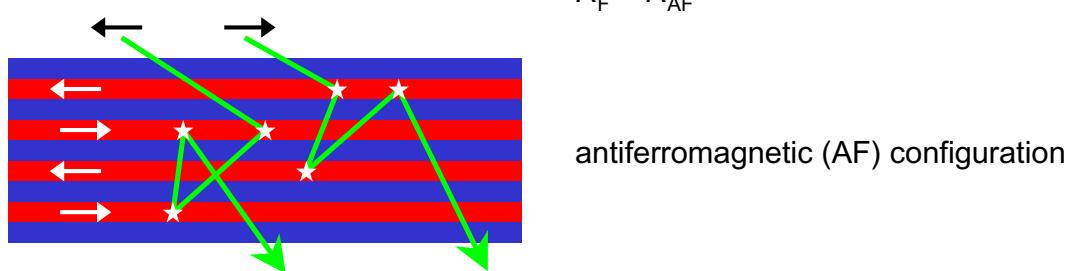
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- mechanism of GMR: spin-dependent scattering

two-current model



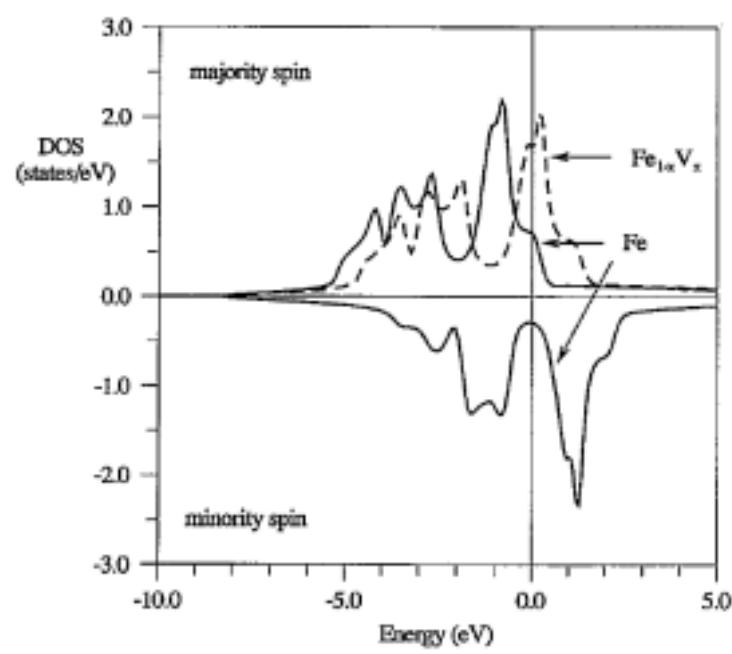
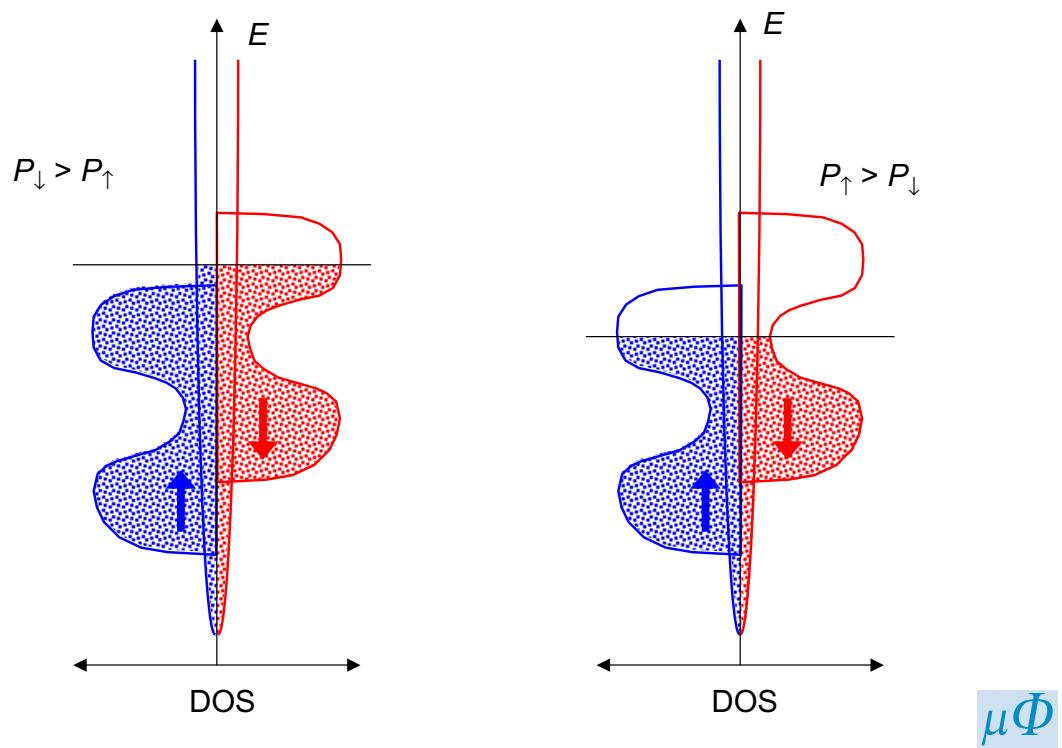
$$R_F < R_{AF}$$



$$A \equiv \frac{R_{AF} - R_F}{R_{AF} + R_F} \quad \text{can be larger than 50%}$$

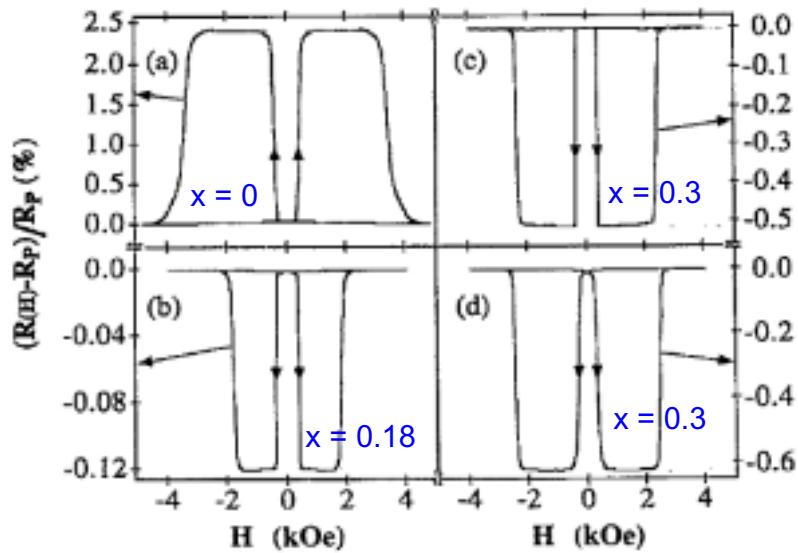
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spin-dependent scattering probability



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$\text{Fe}_{1-x}\text{V}_x / \text{Au} / \text{Co}$



Renard et al., PRB 51, 12821 (1995)

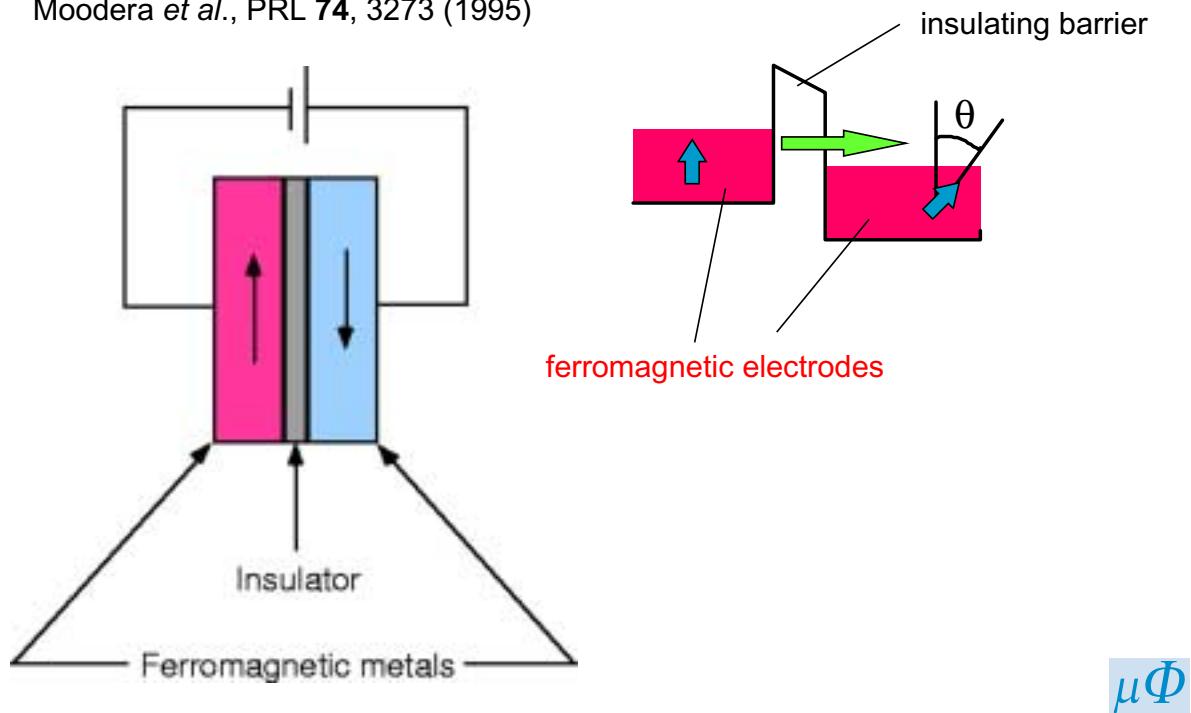
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tunneling-magnetoresistance

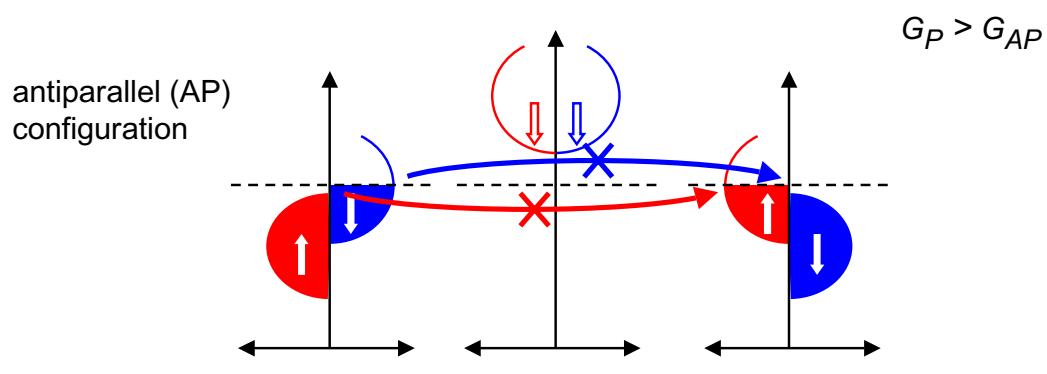
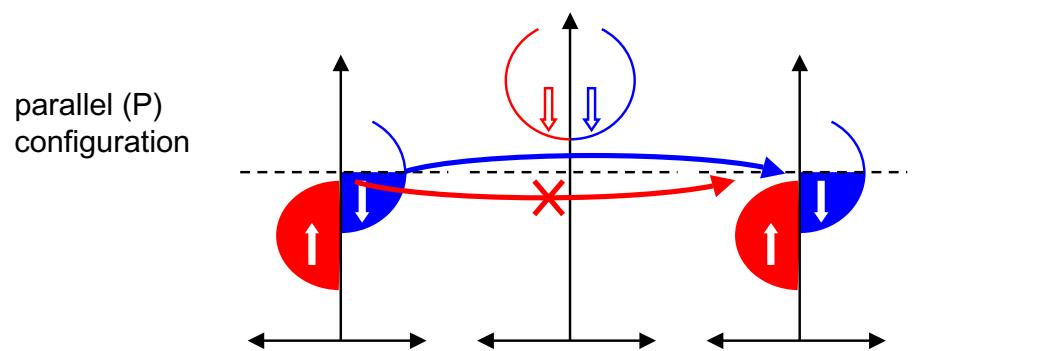
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Tunneling magneto-resistance (TMR)

Jullière, Phys. Lett. **54A**, 225 (1975)
 Moodera *et al.*, PRL **74**, 3273 (1995)

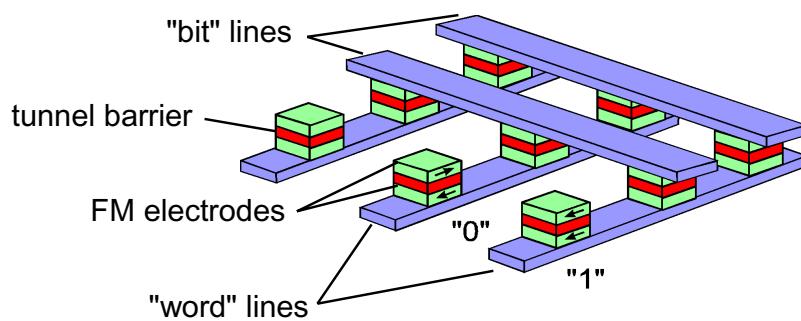


Mechanism of tunneling magneto-resistance



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Applications of TMR: magnetic random access memories (M-RAM)

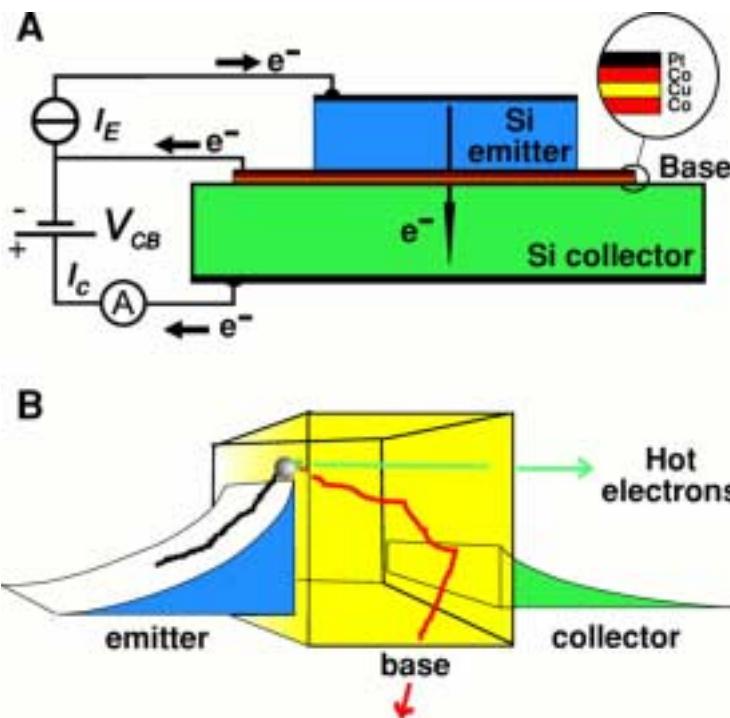


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hot-electron spin-transistor

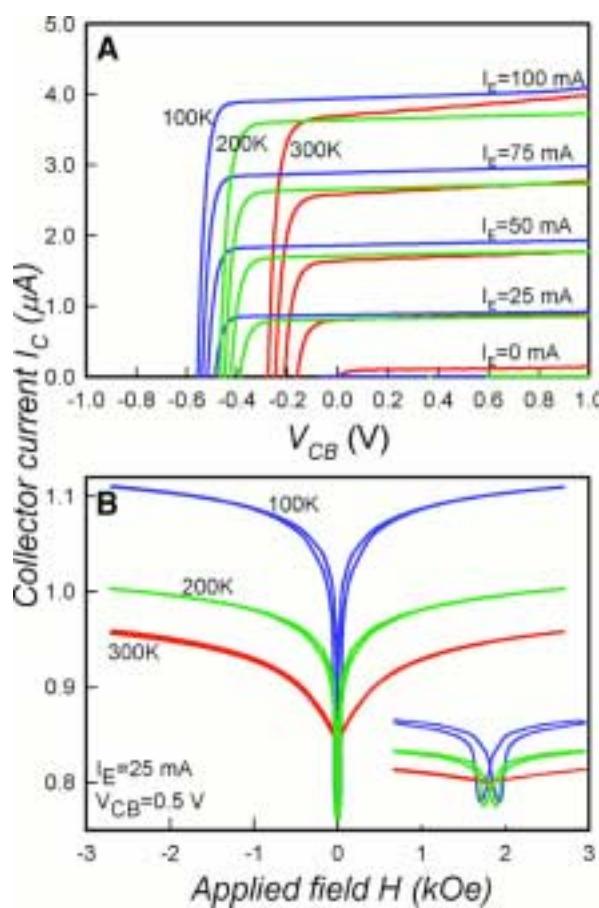
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hot-electron spin-transistor



Monsma et al. PRL 74, 5260 (1995)
Science 281, 407 (1998)

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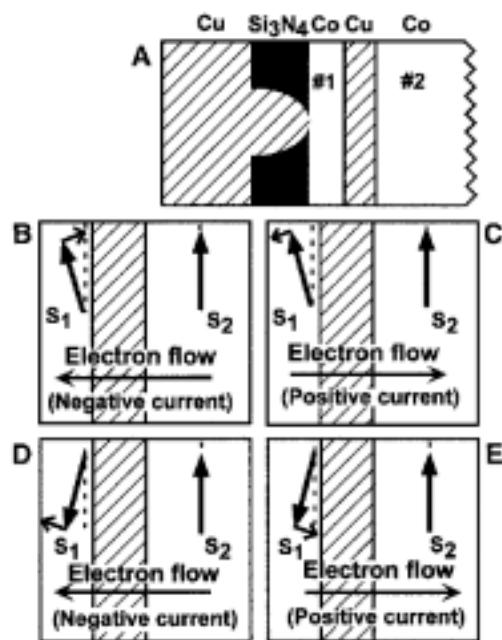


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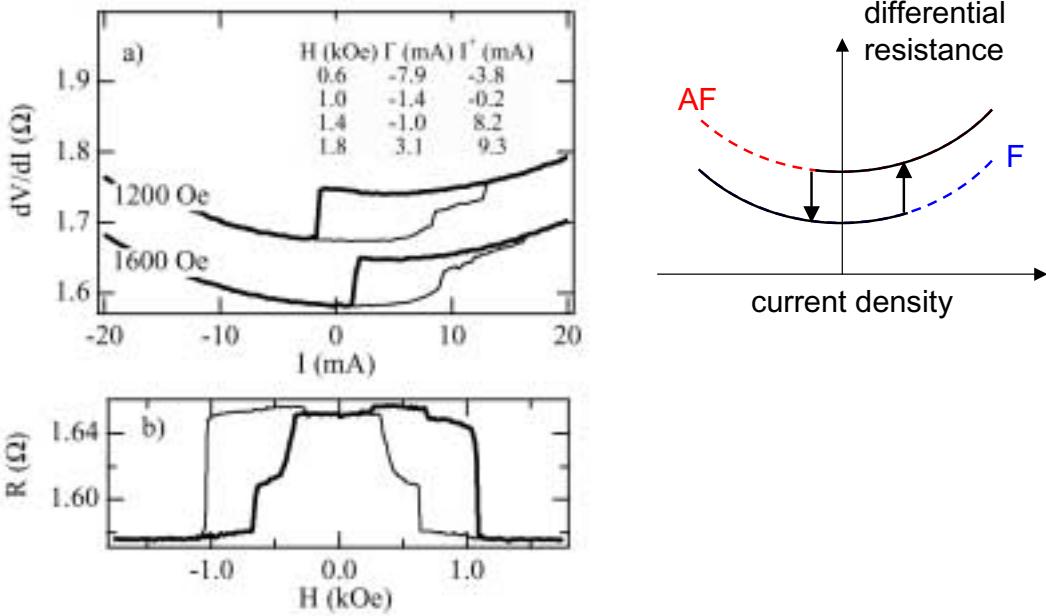
magnetization switching due to spin-injection

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Magnetization switching due to spin-injection



$\mu\Phi$



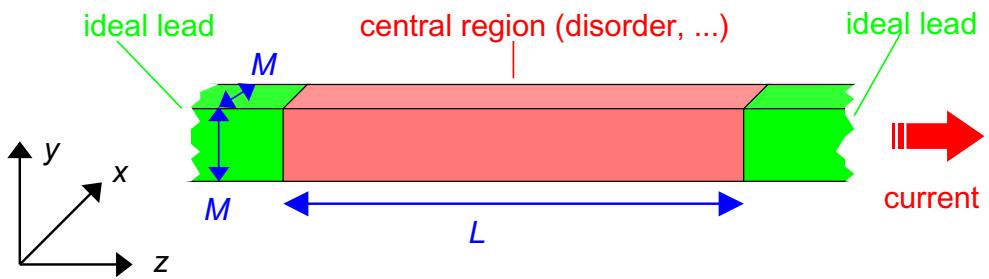
Myers *et al.*, Science **285**, 867 (1999)
 Katine *et al.*, PRL 84, 3149 (2000)

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ab initio calculations of
 perpendicular current GMR

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model considered:



periodic repetition of the supercell in x and y directions

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Conductances within the Landauer-Büttiker formalism

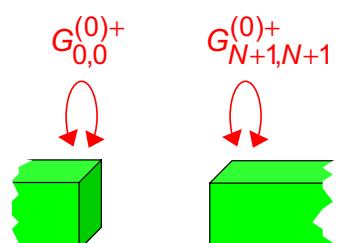
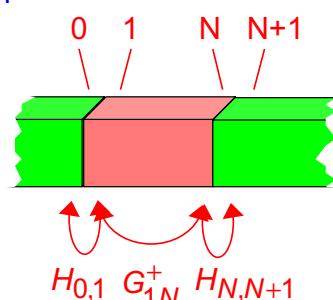
$$C = C^\uparrow + C^\downarrow \quad (\text{spin-colinear case})$$

$$C^\sigma = \frac{e^2}{h} \frac{1}{N_{||}} \sum_{\mathbf{k}_{||}} T^\sigma(\mathbf{k}_{||}, \varepsilon_F)$$

transmittance for spin σ , wavevector $\mathbf{k}_{||}$ and energy ε_F :

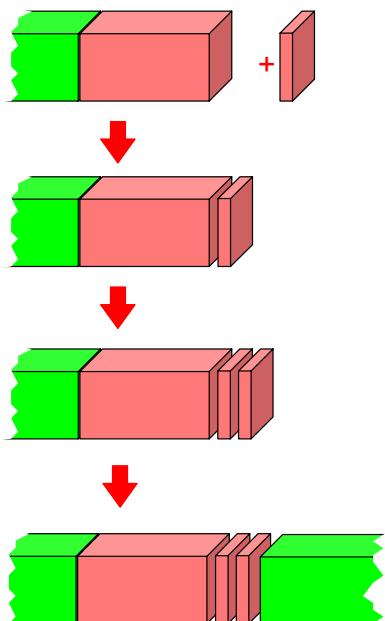
$$T^\sigma(\mathbf{k}_{||}, \varepsilon_F) = \text{Tr} [\Gamma_1 G_{1,N}^+ \Gamma_N G_{N,1}^-]$$

$$\Gamma_1 = i H_{1,0} (G_{0,0}^{(0)+} - G_{0,0}^{(0)-}) H_{0,1}$$

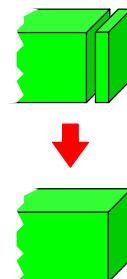


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recursive calculation of
the Green's function



calculation of the surface Green's function:
layer addition (or removal) invariance
→ self-consistent (Dyson-like) equation



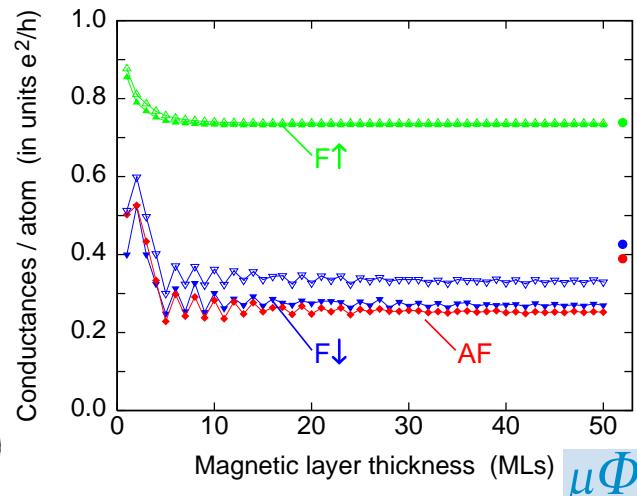
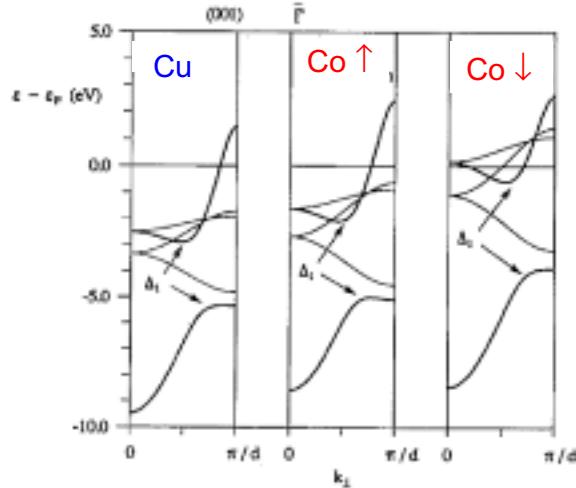
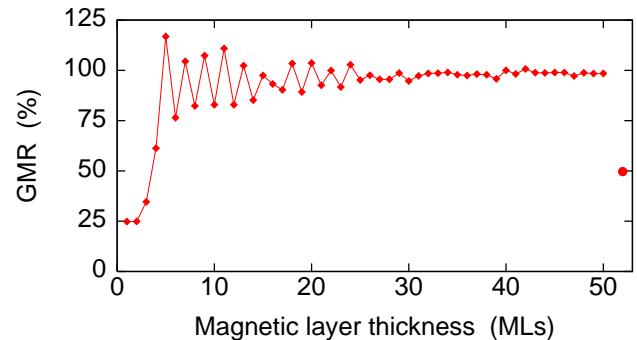
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Computational details:

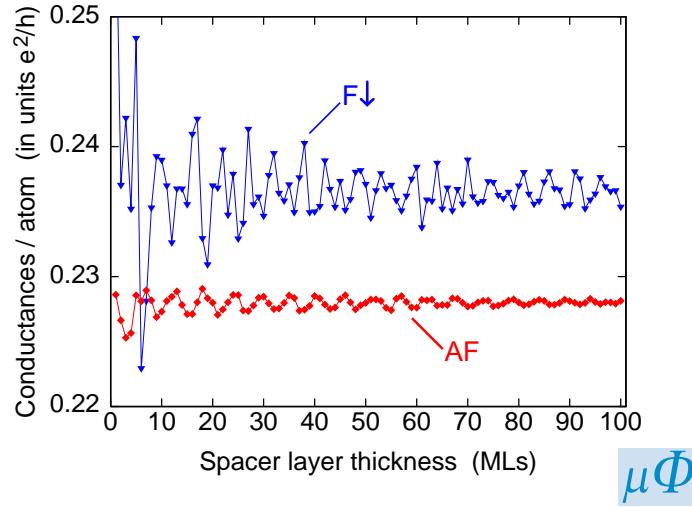
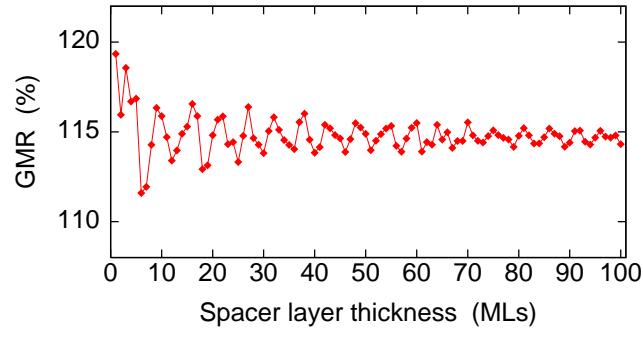
- density functional theory (local density approximation)
- TB-LMTO method (Green's function)
well adapted to surface problems
- imaginary energy for GF calculations: 10^{-7} Ry
- disordered systems:
 - 5×5 supercell averaged over 5 configurations, or
 7×7 supercell averaged over 3 configurations
 - on-site potential parameters obtained from (layer dependent) CPA method
- k_{\parallel} -integration: 10 000 points in the full fcc(001) SBZ
- definition of GMR ratio: $GMR \equiv \frac{C_F - C_{AF}}{C_{AF}}$

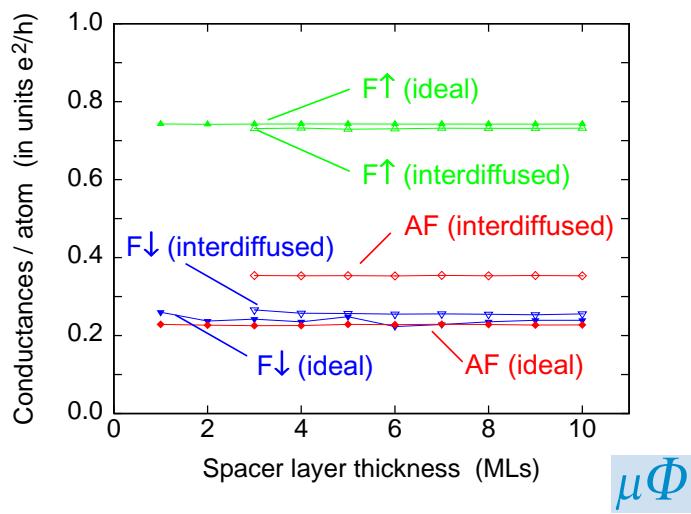
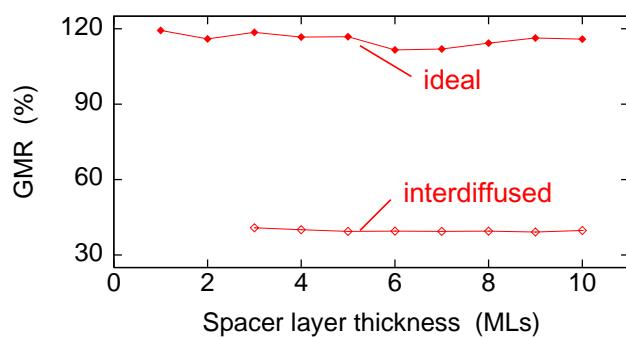
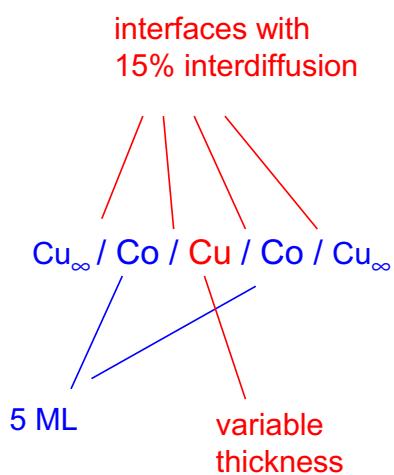
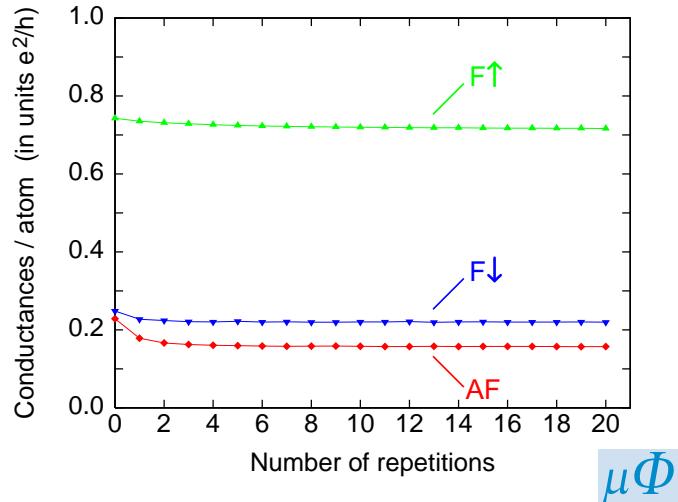
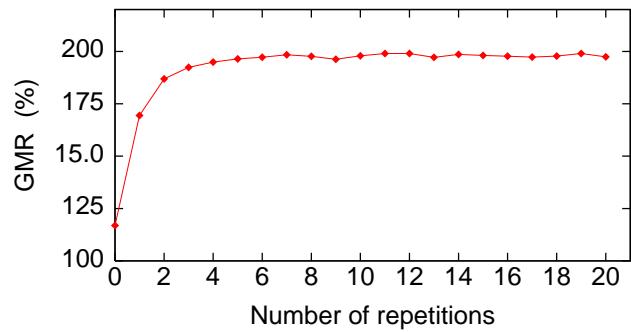
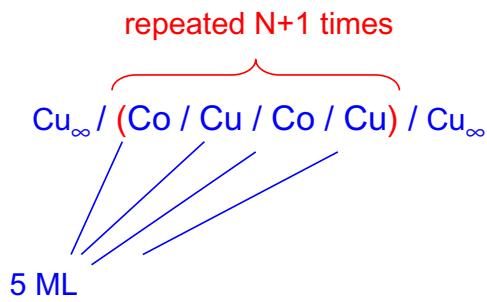
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$\text{Cu}_\infty / \text{Co} / \text{Cu} / \text{Co} / \text{Cu}_\infty$
variable thickness
5 ML

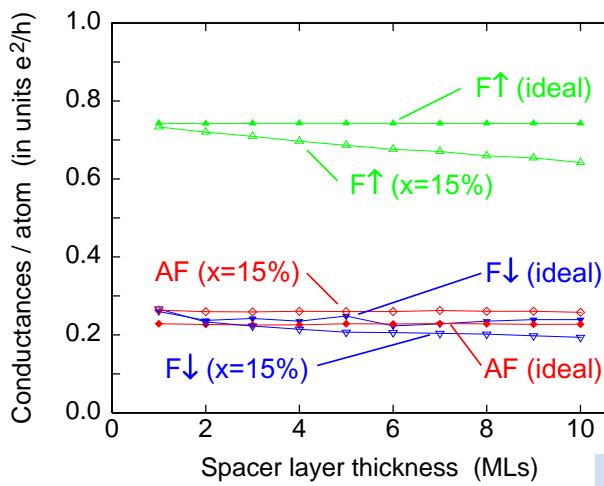
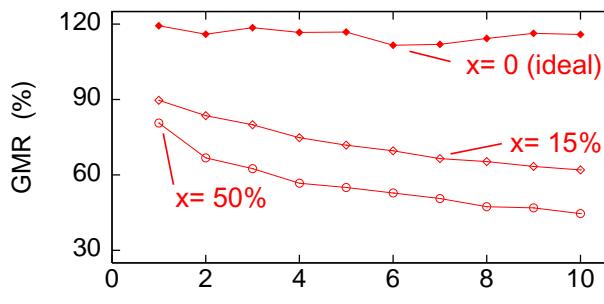


$\text{Cu}_\infty / \text{Co} / \text{Cu} / \text{Co} / \text{Cu}_\infty$
5 ML
variable thickness



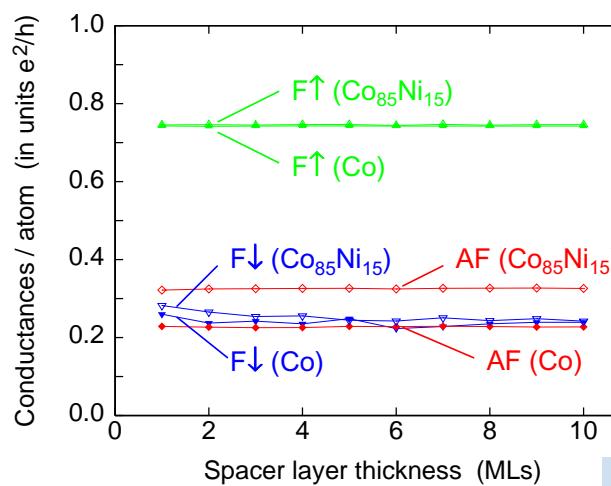
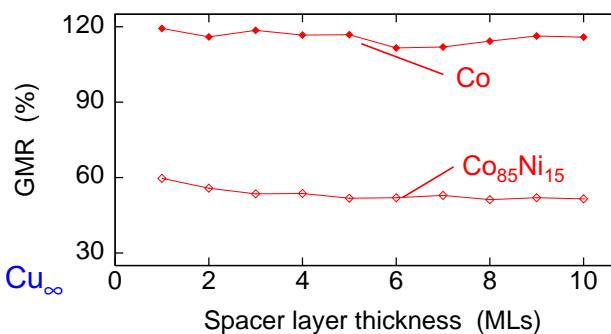


$\text{Cu}_\infty / \text{Co} / \text{Cu}_{(1-x)}\text{Pd}_x / \text{Co} / \text{Cu}_\infty$
 5 ML
 variable thickness



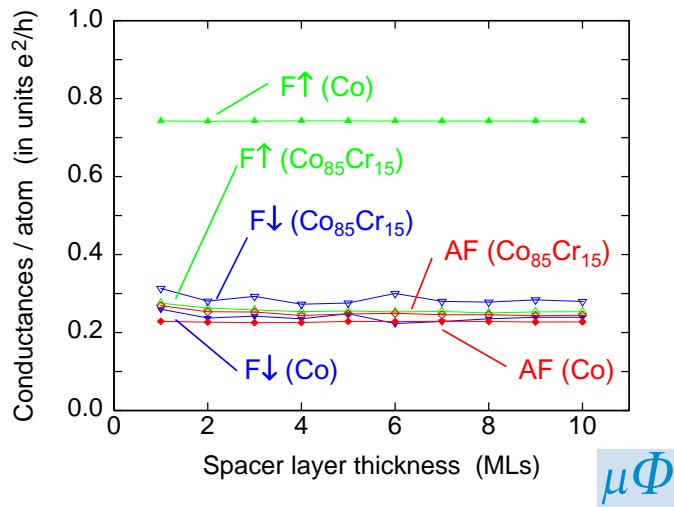
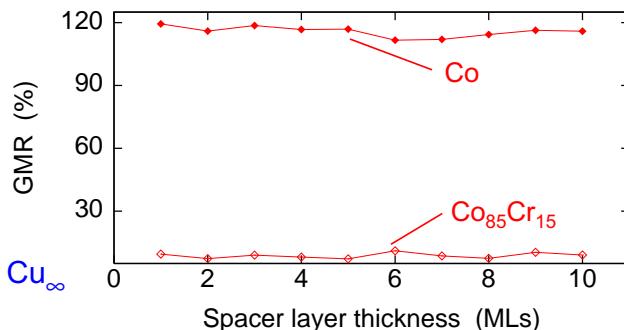
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$\text{Cu}_\infty / \text{Co}_{(1-x)}\text{Ni}_x / \text{Cu} / \text{Co}_{(1-x)}\text{Ni}_x / \text{Cu}_\infty$
 5 ML
 variable thickness



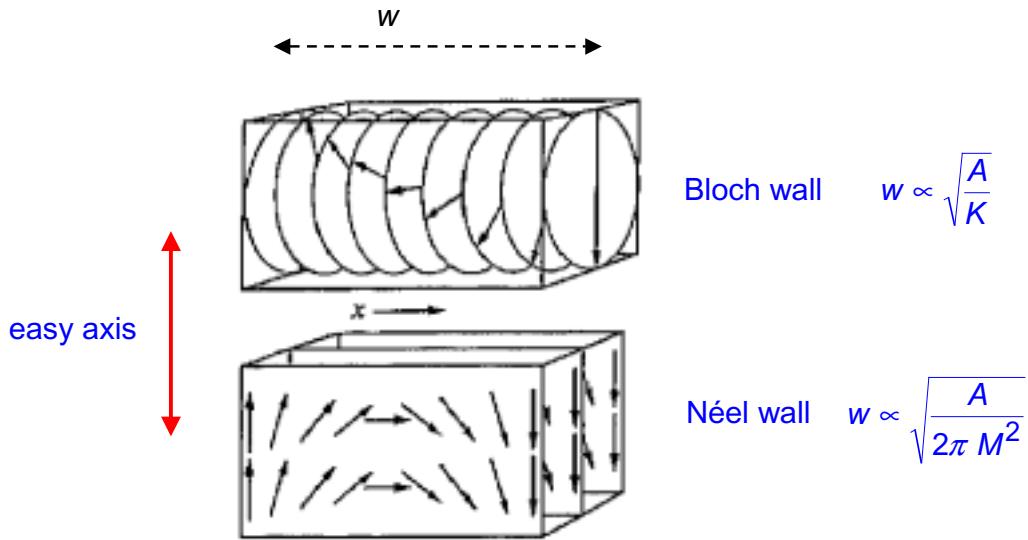
$\mu\Phi$

$\text{Cu}_\infty / \text{Co}_{(1-x)}\text{Cr}_x / \text{Cu} / \text{Co}_{(1-x)}\text{Cr}_x / \text{Cu}_\infty$
 5 ML variable thickness

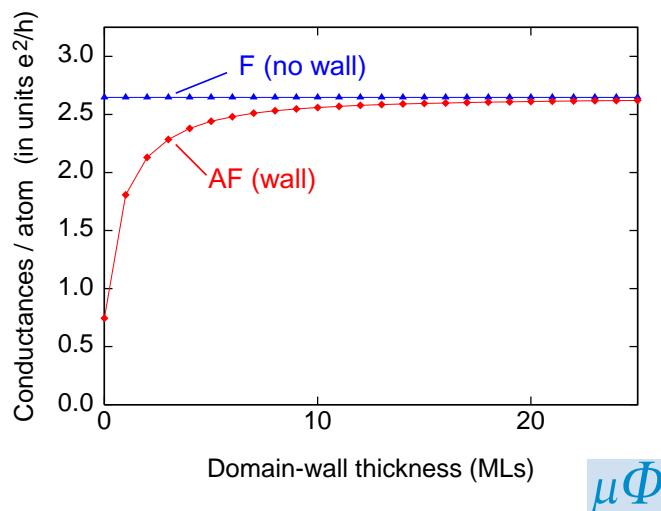
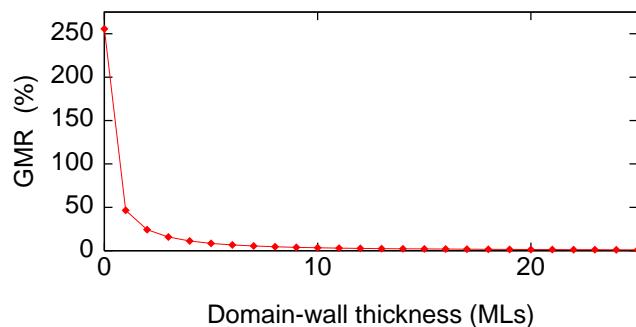


domain wall magnetoresistance

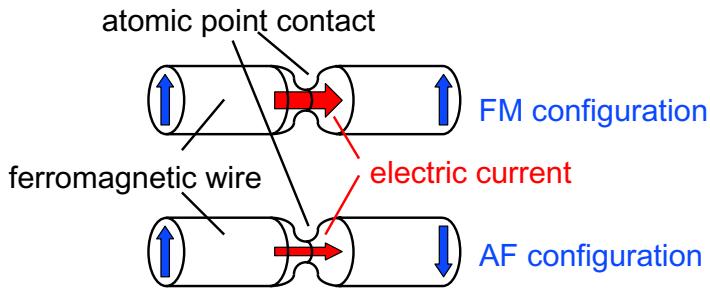
$\mu\Phi$



$\text{Co}_\infty / \text{Co} / \text{Co}_\infty$
 magnetic wall of variable thickness
 (linear rotation of magnetization)



Giant magnetoresistance of ferromagnetic atomic point contacts



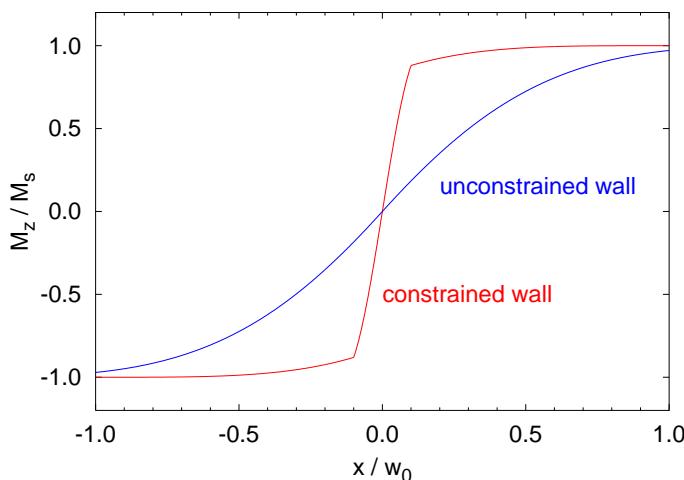
explanation requires:

- narrow magnetic wall in an atomic point contact
- large resistance due to a narrow wall

$$\mu\Phi$$

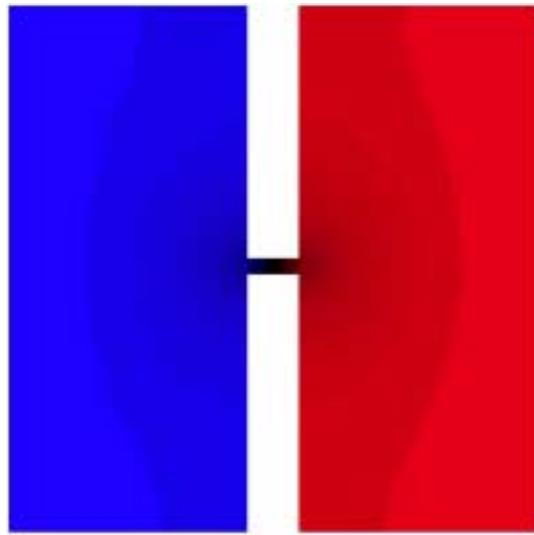
geometrical constriction → new kind of magnetic wall:

- structure (almost) entirely determined by the constriction geometry
- energy = (almost) pure exchange energy
- width determined by the characteristic size of the constriction
→ wall can be extremely narrow in an atomic point contact

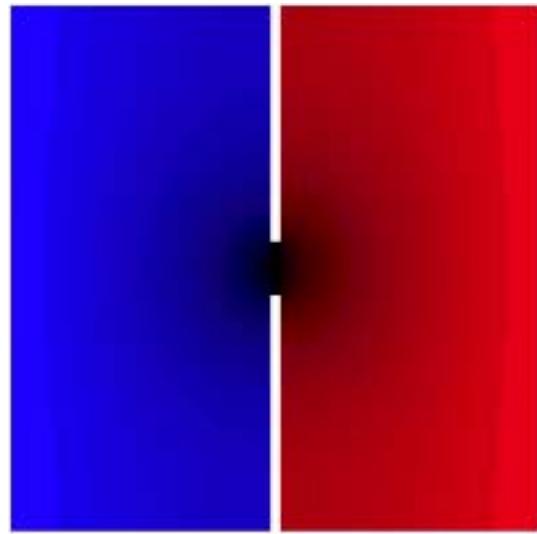


P. Bruno, PRL 83, 2425 (1999)

$$\mu\Phi$$



20 nm X 20 nm

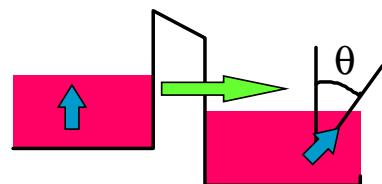


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**theory of tunneling
magnetoresistance**

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simple approach: Jullière model

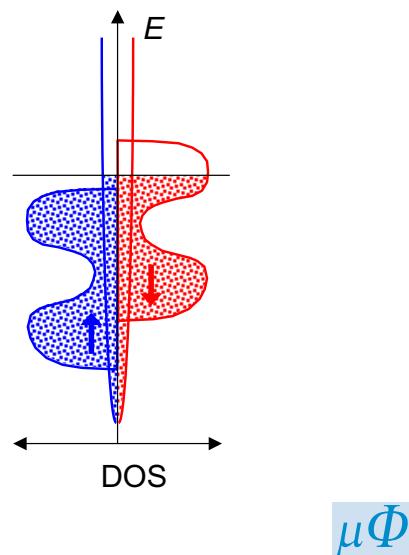


assumptions: $G = G_{\uparrow} + G_{\downarrow}$

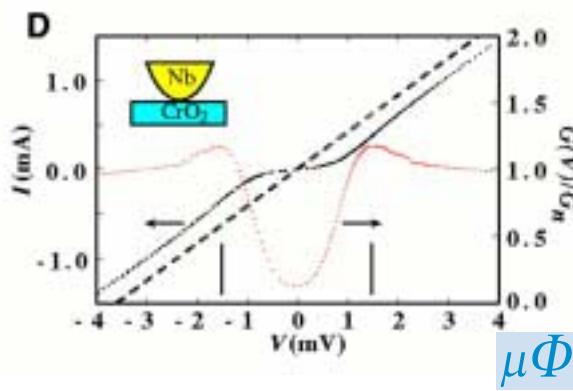
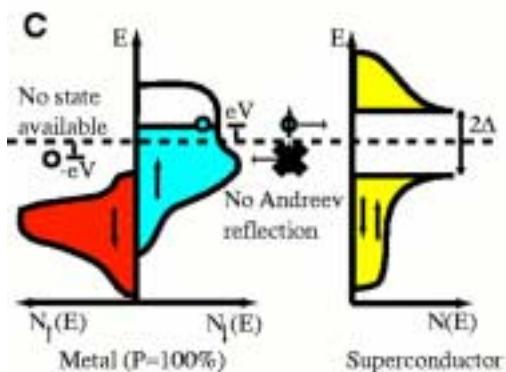
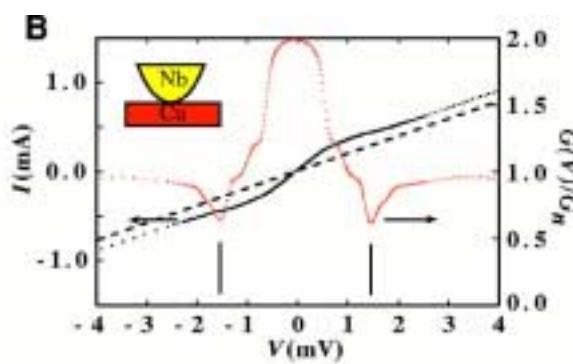
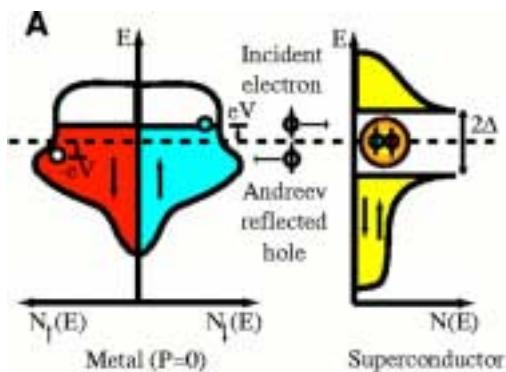
$$G\sigma \propto \rho_1^{\sigma}(\varepsilon_F) \rho_2^{\sigma}(\varepsilon_F)$$

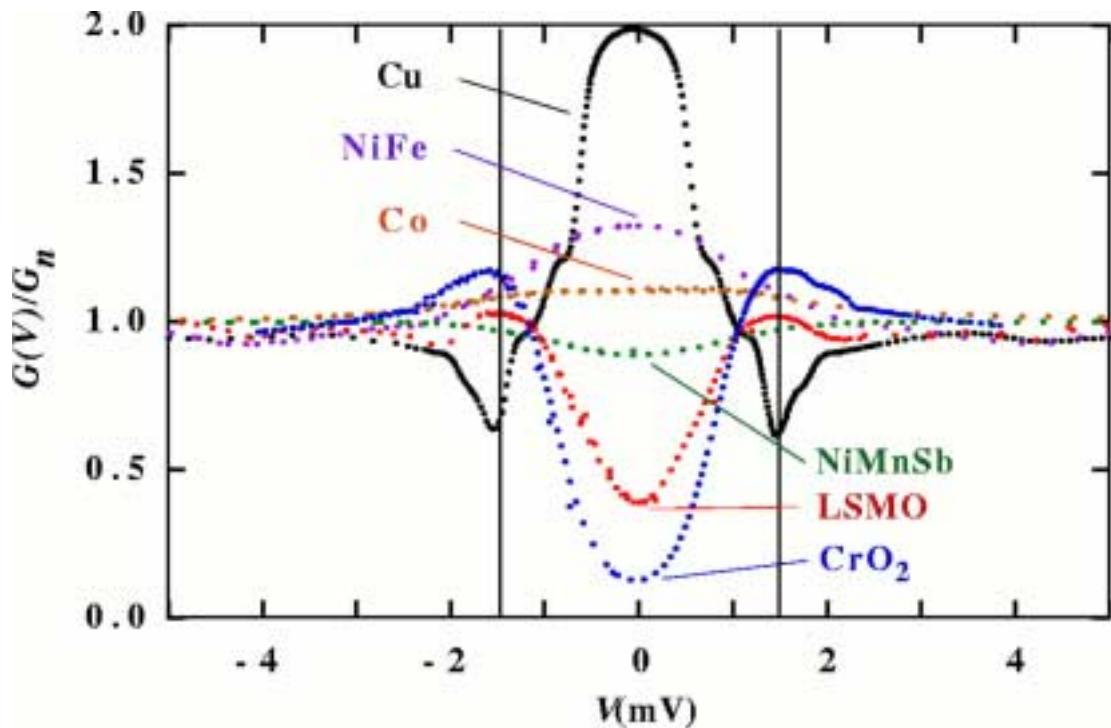
$$\rightarrow A \equiv \frac{G_{AP} - G_P}{G_{AP} + G_P} = P_1 P_2$$

$$\text{with } P \equiv \frac{\rho_{\uparrow}(\varepsilon_F) - \rho_{\downarrow}(\varepsilon_F)}{\rho_{\uparrow}(\varepsilon_F) + \rho_{\downarrow}(\varepsilon_F)}$$



measurement of the spin-polarization by spin-injection into a superconductor

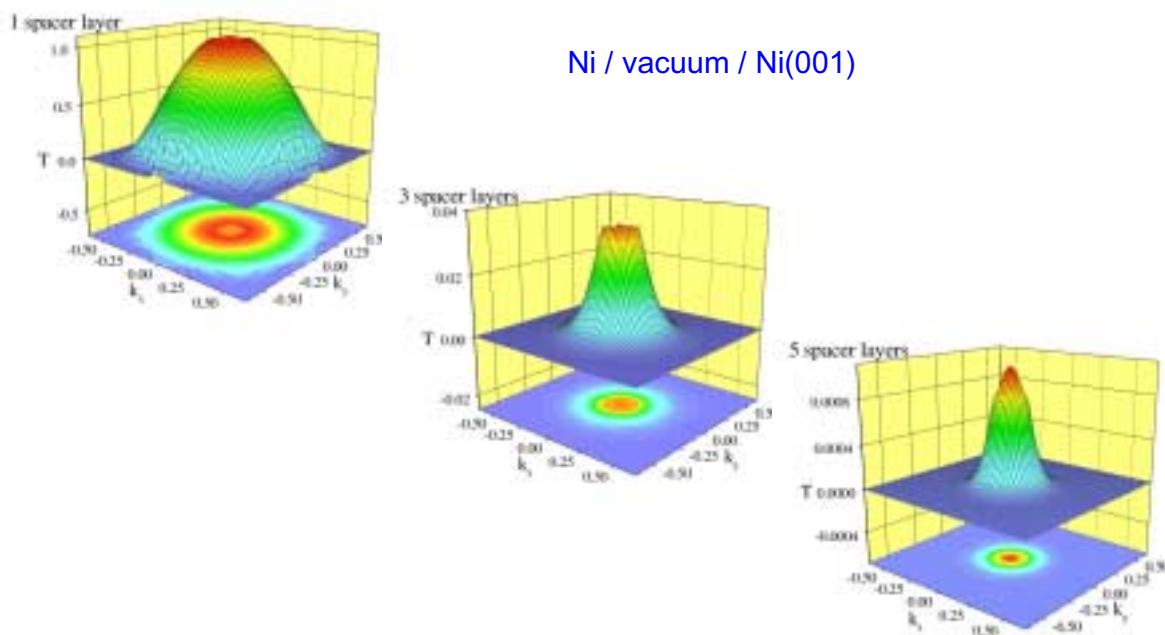




Soulen *et al.*, Science **282**, 85 (1998)

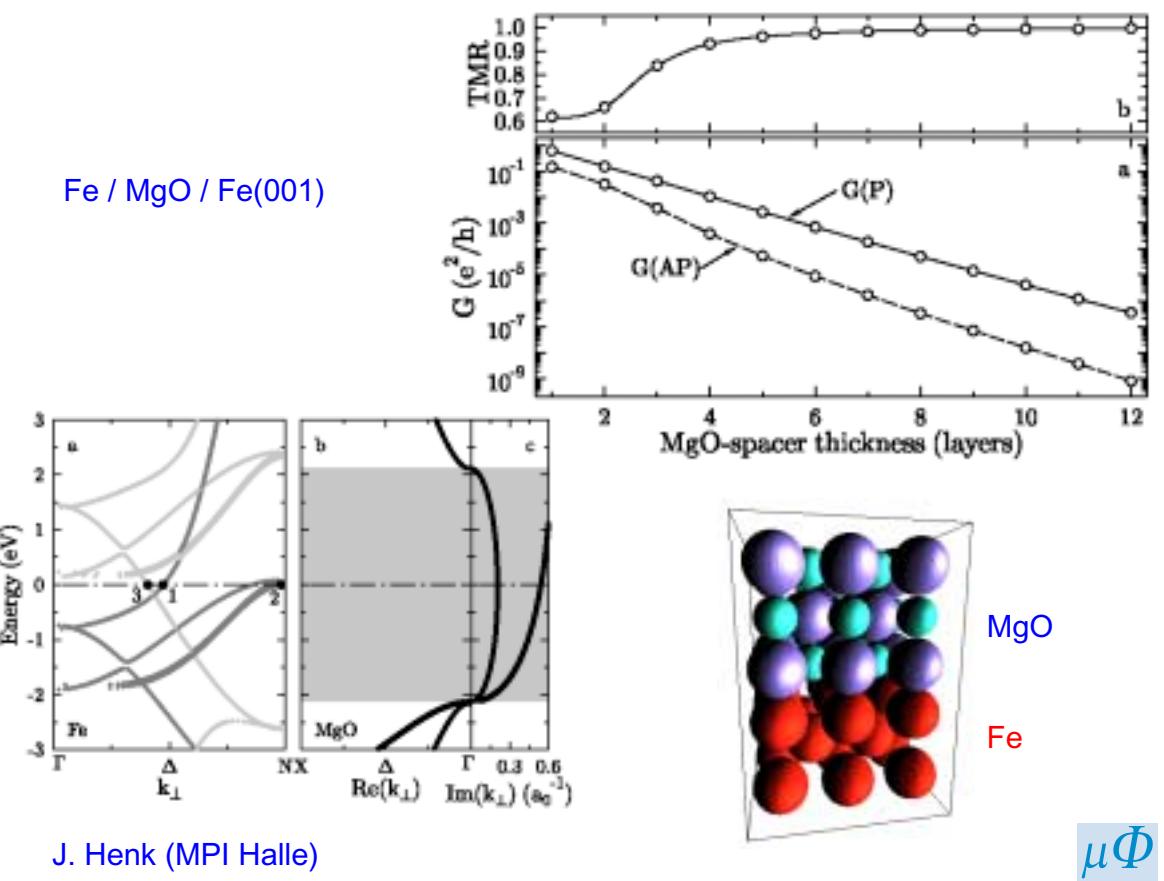
$\mu\Phi$

ab initio calculation of tunnel conductance



J. Henk (MPI Halle)

$\mu\Phi$



J. Henk (MPI Halle)